

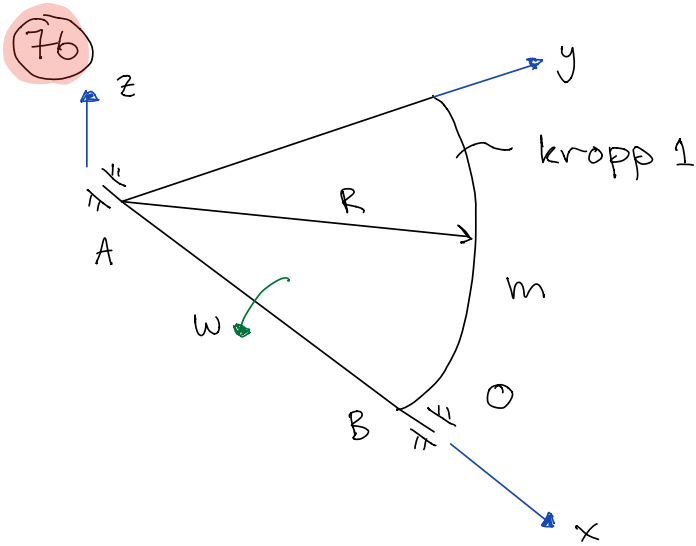
Föreläsning 14

TMME04 – Mekanik II

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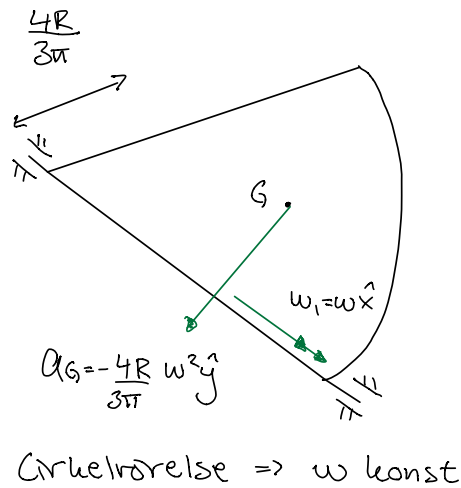
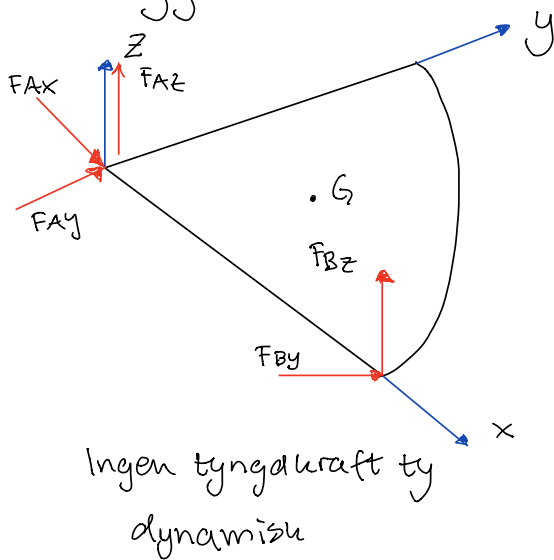
<https://www.instagram.com/olwettergren/>



Givet: w konst, $Axyz$ kroppsfix

Sökt: Dynamiska reaktionskrafter \bar{F}_A och \bar{F}_B

Frilägg:



Euler I:

$$\bar{F} = m \bar{a}_G$$

$$\hat{x}: F_{Ax} = 0 \quad (1)$$

$$\hat{y}: F_{Ay} + F_{By} = m a_{gy} = m \left(-\frac{4R}{3n} \omega^2 \right) \quad (2)$$

$$\hat{z}: F_{Az} + F_{Bz} = 0 \quad (3)$$

Euler II:

$$\bar{M}_A = \dot{\bar{h}}_A, \quad A \text{ fix i i-ram.}$$

$$\bar{h}_A = I_A \bar{\omega}, \quad A \text{ även fix i kropp 1}$$

$$\bar{M}_A = \frac{d\bar{h}_A}{dt} + \bar{\omega}_r \times \bar{h}_A \quad (4)$$

Val av referensram r att derivera i:

$$r = 1$$

Inför $Axyz$ fixt i kropp 1

$$\bar{h}_A = \begin{bmatrix} I_{Axx} & * & * \\ I_{Axy} & * & * \\ 0 & * & * \end{bmatrix} \begin{bmatrix} \omega \\ 0 \\ 0 \end{bmatrix} = I_{Axx} \omega \hat{x} + I_{Axy} \omega \hat{y}$$

$$I_{Axz} = \int_{\substack{V \\ z=0}} -xz \, dm$$

$$\left(\frac{d\bar{h}_A}{dt} \right)_1 = \bar{0}$$

Insättning i (4) \Rightarrow

$$\bar{M}_A = \omega \hat{x} \times (I_{Axx} \omega \hat{x} + I_{Axy} \omega \hat{y}) = I_{Axy} \omega^2 \hat{z} \quad (5)$$

Men \bar{M}_A i v.l ges av friläggingsfiguren

$$M_{Ax} = \overset{(5)}{0} = 0$$

$$M_{Ay} = -F_{Bz} R \stackrel{(5)}{=} 0 \quad (6)$$

$$M_{Az} = F_{By} R \stackrel{(5)}{=} I_{Axy} \omega^2 \quad (7)$$

(1) \Rightarrow

$$F_{Ax} = 0$$

(6), (3) \Rightarrow

$$F_{Az} = F_{Bz} = 0$$

(7) \Rightarrow

$$F_{By} = \frac{I_{Axy} \omega^2}{R} \quad (8)$$

Insättning i (2) \Rightarrow

$$F_{Ay} = -\frac{I_{Axy} \omega^2}{R} - \frac{4mR}{8\pi} \omega^2 \quad (9)$$

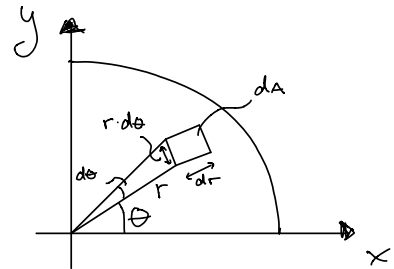
I_{Axy} ?

$$I_{Axy} = -\int xy \, dm =$$

$$= \left/ \begin{array}{l} x = r \cos \theta, \quad y = r \sin \theta \\ dm = m \cdot dA = \frac{4m}{\pi R^2} dr \cdot r d\theta \end{array} \right/ =$$

$$= -\frac{4m}{\pi R^2} \int_0^{\pi/2} \int_0^R r^2 \sin \theta \cos \theta \, dr d\theta = -\frac{2m}{\pi R^2} \cdot \frac{1}{4} R^2 \int_0^{\pi/2} \sin 2\theta \, d\theta =$$

$$= -\frac{mR^2}{4\pi} \left[-\cos 2\theta \right]_0^{\pi/2} = \frac{-mR^2}{2\pi}$$



(8) \Rightarrow

$$F_{By} = - \frac{mR}{2\pi} \omega^2$$

(9) \Rightarrow

$$F_{Ay} = - \frac{5mR}{6\pi} \omega^2$$

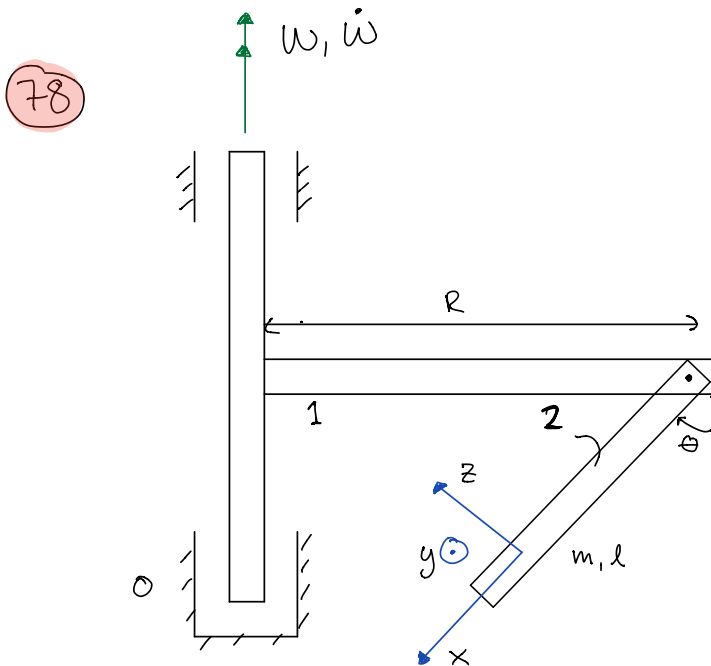
$$\therefore \bar{F}_B = - \frac{mR}{2\pi} \omega^2 \hat{y}$$

$$\bar{F}_A = - \frac{5mR}{6\pi} \omega^2 \hat{y}$$

Ser att $F=0$ då $\omega=0$, så för alla dynamiska krafter.

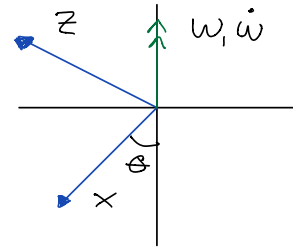
Dessutom kraften vid $A > B$.

Standard: Euler I & II. Tröghetsprodukten skiljer



Givet: G_{xyz} fixt i stängen

Sökt: $\bar{\alpha}_2, T, \bar{h}_0$



a)

$$\bar{w}_2 = w_{2/0} = \bar{w}_{2/1} + \bar{w}_{1/0} = -\dot{\theta} \hat{y} + w(-\cos\theta \hat{x} + \sin\theta \hat{z}) \quad (1)$$

$$\begin{aligned} \bar{\alpha}_2 &= \bar{\alpha}_{2/0} = \left(\frac{d\bar{w}_{2/0}}{dt} \right)_0 = \left(\frac{d\bar{w}_{2/1}}{dt} \right)_0 + \underbrace{\left(\frac{d\bar{w}_{1/0}}{dt} \right)_0}_{\bar{\alpha}_{1/0}} = \\ &= \underbrace{\left(\frac{d\bar{w}_{2/1}}{dt} \right)}_{\bar{\alpha}_{2/1}} + \bar{w}_{1/0} \times \bar{w}_{2/1} + \bar{\alpha}_{1/0} = \end{aligned}$$

$$\begin{aligned} &= -\ddot{\theta} \hat{y} + w(-\cos\theta \hat{x} + \sin\theta \hat{z}) \times (-\dot{\theta} \hat{y}) + \\ &+ \dot{w}(-\cos\theta \hat{x} + \sin\theta \hat{z}) = (w\dot{\theta}\sin\theta - \dot{w}\cos\theta) \hat{x} - \\ &-\dot{\theta} \hat{y} + (w\dot{\theta}\cos\theta + \dot{w}\sin\theta) \hat{z} \end{aligned}$$

Alt:

$$\begin{aligned} \bar{\alpha}_{2/0} &= \left(\frac{d\bar{w}_{2/0}}{dt} \right)_0 = \left(\frac{d\bar{w}_{2/0}}{dt} \right)_2 + \underbrace{\bar{w}_{2/0} \times \bar{w}_{2/0}}_{=0} = \\ &= -\ddot{\theta} \hat{y} + (-\dot{w}\cos\theta + w\sin\theta \dot{\theta}) \hat{x} + (\dot{w}\sin\theta + w\cos\theta \dot{\theta}) \hat{z} \end{aligned}$$

b)

$$T = \frac{1}{2} m \bar{v}_G \cdot \bar{v}_G + \frac{1}{2} \bar{\omega}_2 \cdot \bar{h}_G \quad (2)$$

$$\bar{v}_G = \underbrace{\bar{v}_O}_{-R\omega\hat{y}} + \underbrace{\bar{\omega}_2 \times \bar{r}_{OG}}_{l/2\hat{x}} \stackrel{(1)}{=} -R\omega\hat{y} + \frac{l}{2}\dot{\theta}\hat{z} + \frac{l}{2}\omega\sin\theta\hat{y} \quad (3)$$

Circletravel

$\bar{v}_G \neq \bar{v}_A + \bar{\omega}_2 \times \bar{r}_{AG}$ fungerar ej ty ej
samma kropp

$$\bar{h}_G = \bar{I}_G \bar{\omega}_2 = \begin{bmatrix} 0 & 0 & 0 \\ 0 & ml^2/12 & 0 \\ 0 & 0 & ml^2/12 \end{bmatrix} \begin{bmatrix} -\omega\cos\theta \\ -\dot{\theta} \\ \omega\sin\theta \end{bmatrix} =$$

$$= \frac{ml^2}{12}\dot{\theta}\hat{y} + \frac{ml^2}{12}\omega\sin\theta\hat{z} \quad (4)$$

Insättning i (2) \Rightarrow

$$T = \frac{ml^2}{6}\dot{\theta} + \frac{ml^2}{6}\omega^2\sin^2\theta + \frac{mR^2}{2}\omega^2 - \frac{mRl}{2}\omega^2\sin\theta$$

c)

$$\bar{h}_O = \bar{h}_G + \bar{r}_{OG} \times m\bar{v}_G \stackrel{(3),(4)}{=} -\frac{ml^2}{3}\dot{\theta}\hat{y} + mR\omega\left(\frac{l}{3}\sin\theta - \frac{R}{2}\right)\hat{z}$$